

Optimal Perturbations with the MITgcm: MOC & tropical SST in an idealized ocean

Laure Zanna (Harvard)

Patrick Heimbach, Eli Tziperman, Andy Moore

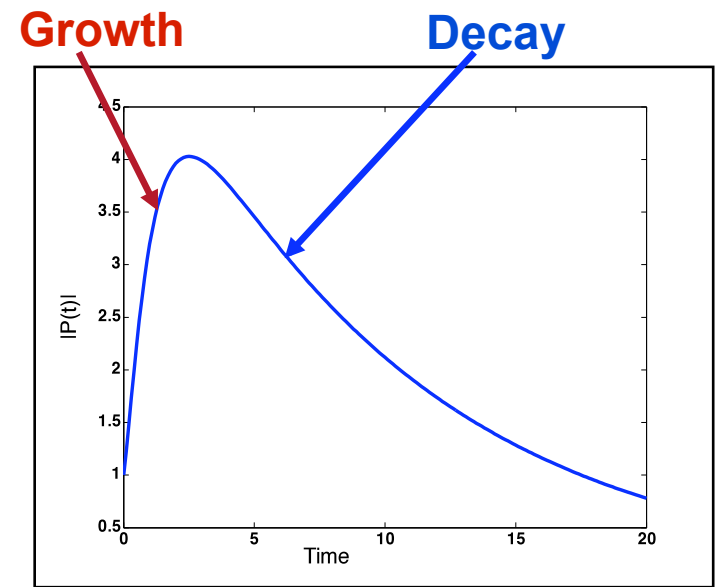
ECCO2 Meeting, Sep 23rd 2008

Which perturbations can lead to the most efficient growth?

- In stable systems, perturbations can **grow** significantly before eventually decaying due to the interaction of several non-orthogonal modes (e.g., Farrell 1988, Trefethen 1993)

- **Optimal initial conditions** = singular vectors → fastest growing perturbations leading to an amplification of a given quantity (e.g., Buizza & Palmer 1995)

- **Relevance** = e.g., climate stability & variability, sensitivity, predictability and error growth, building an observational system (e.g., Marotzke et al 1999, Moore & Kleeman 1999, Moore et al 2004)



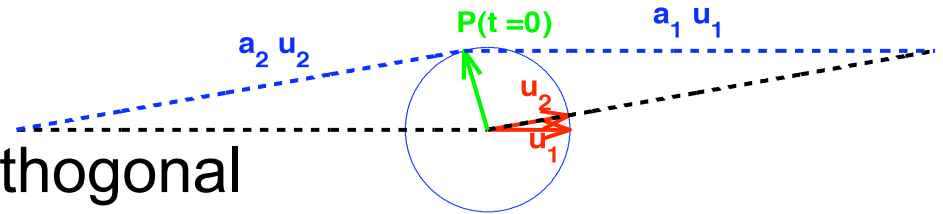
Objectives

- **Spatial structure of the optimal initial conditions** leading to the maximum growth of the physical quantities: heat flux, MOC, tropical SST, kinetic & available potential energy, ...
- Identification of the **growth mechanism** for the perturbations & implications for stability and variability of ocean & climate
- → Can observed ocean variability be explained as small amplitude **damped linear dynamics excited by atmospheric & other stochastic forcing via non-normal growth?**

Transient Amplification

Stable linear system $\frac{d\vec{P}(t)}{dt} = A\vec{P}(t)$, $\vec{P}(t) \rightarrow 0$ as $t \rightarrow \infty$

If A is non-normal $AA^T \neq A^T A$
 then eigenvectors \vec{u}_i are not orthogonal
 → may lead to **transient amplification**



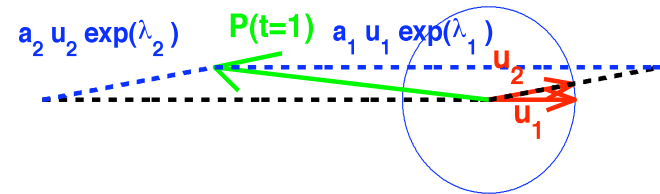
(2D) solution at time τ :

$$\vec{P}(\tau) = a_1 \vec{u}_1 e^{\lambda_1 \tau} + a_2 \vec{u}_2 e^{\lambda_2 \tau}$$

If $\lambda_2 \ll \lambda_1 < 0$, then $a_2 \vec{u}_2 e^{\lambda_2 \tau} \rightarrow 0$ quickly

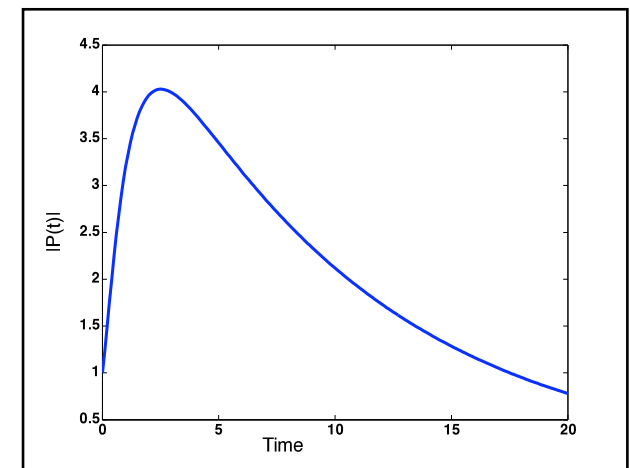
leaving mostly $\vec{P}(t=1) \approx a_1 \vec{u}_1 e^{\lambda_1}$

eventually $\vec{P}(t \rightarrow \infty) \rightarrow 0$



Transient amplification: **Interaction of non orthogonal eigenmodes** b/c of

- (1) Partial initial cancellation
- (2) Different decay rates



(e.g. Farrell, 1988, Trefethen, 1993)

Evaluating the Optimal Initial Conditions: Eigenvalue Problem

- Full nonlinear model linearized about steady state

$$\frac{d\vec{P}'}{dt} = A|_{\bar{P}} \vec{P}', \quad \vec{P}'(t) = e^{At} \vec{P}'_0 = B(t) \vec{P}'_0 \quad \text{most fluid dynamical systems are non-normal}$$

- Maximize MOC or SST anomalies at time $t = \tau$ to find optimal initial conditions \vec{P}'_0

$$\max_{\vec{P}'_0} \left\{ \vec{P}'_0^T B^T X B \vec{P}'_0 - \lambda \left(\vec{P}'_0^T Y \vec{P}'_0 - 1 \right) \right\}$$

- Equivalent to a generalized eigenproblem for optimal initial conditions \vec{P}'_0

$$B^T X B \vec{P}'_0 = \lambda Y \vec{P}'_0$$

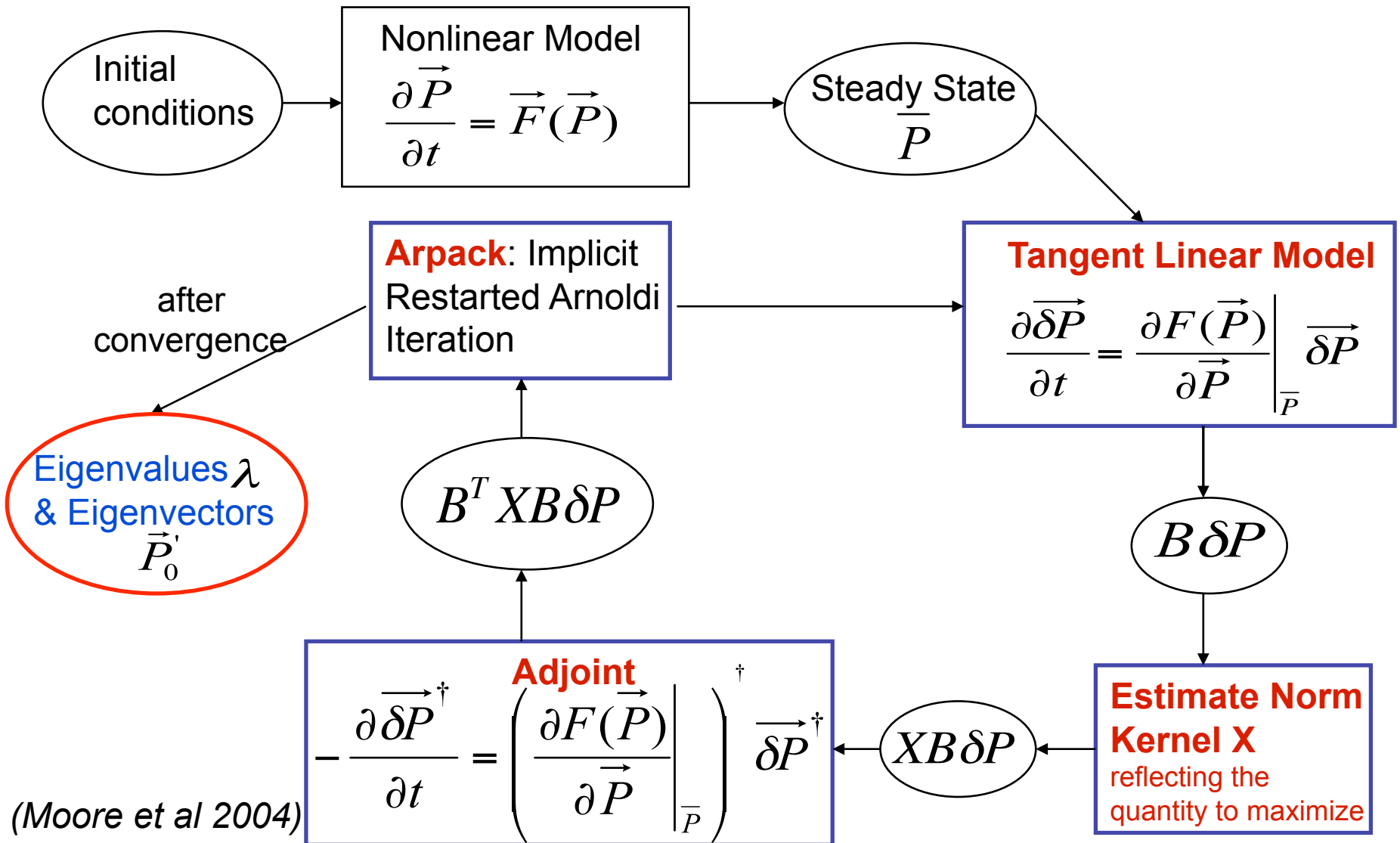
MOC or Tropical
SSTs at $t = \tau$

T and S anomalies at
 $t = 0$

(e.g. Farrell, 1988)

Methodology: Optimals using the MITgcm

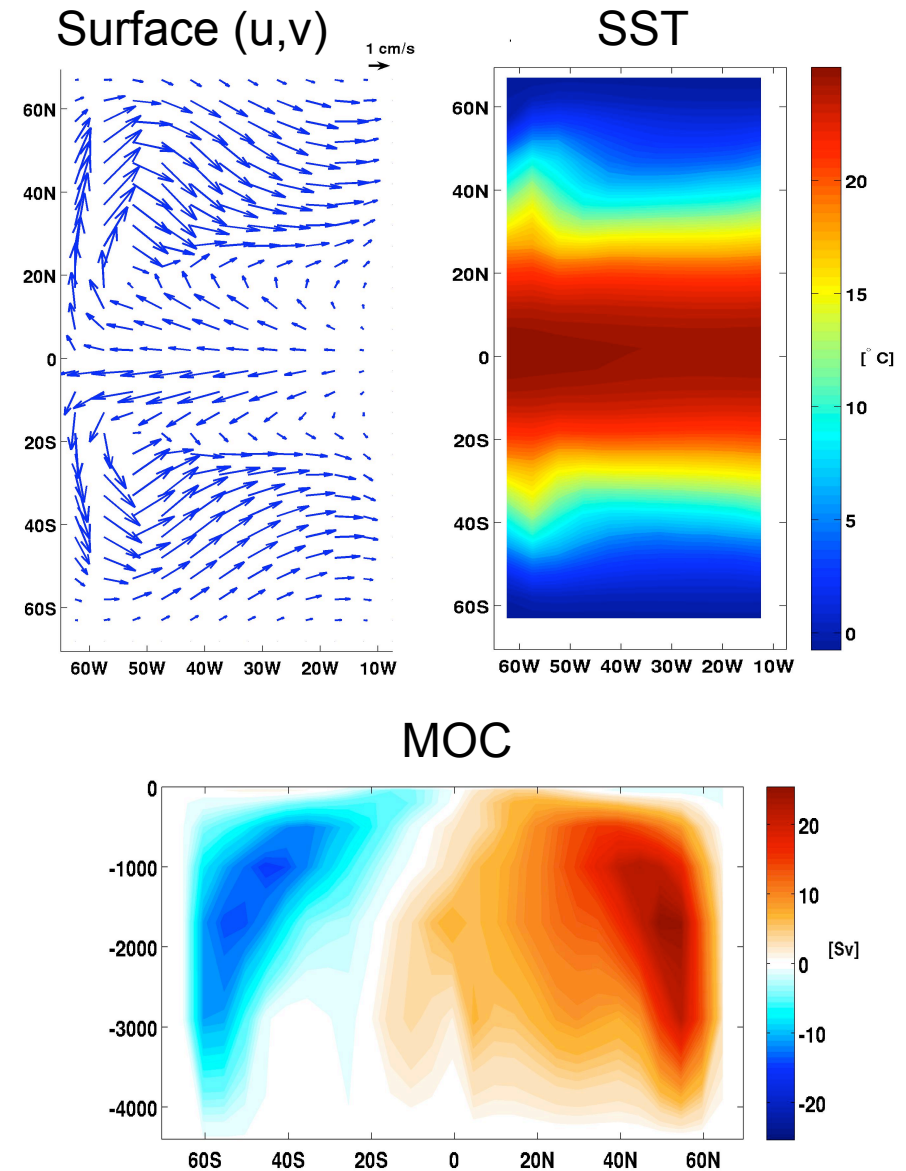
Finding optimal initial conditions \rightarrow Solving for eigenvectors \vec{P}_0' & eigenvalues λ of the generalized eigenproblem $(e^{A\tau})^T X e^{A\tau} (= B^T X B)$



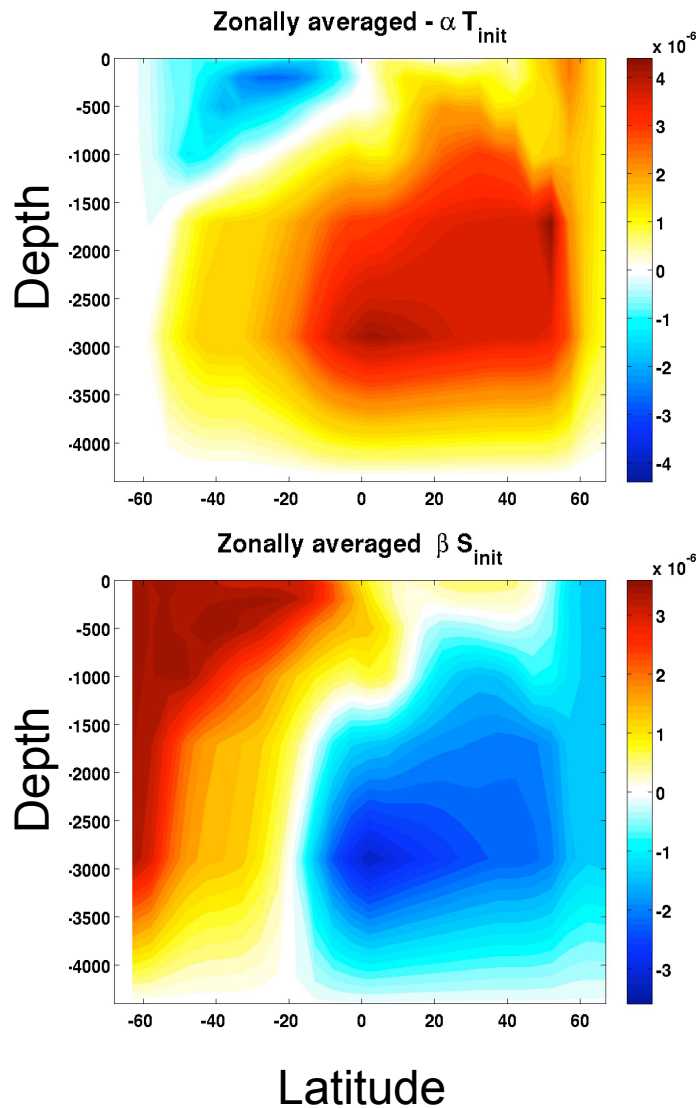
MITgcm: Mean State

- Primitive, hydrostatic, incompressible, Boussinesq eqns on a sphere
- Configuration: rectangular double-hemisphere ocean basin, coarse resolution $3^\circ \times 3^\circ$, 15 vertical levels, flat topography
- Convection=Implicit diffusion
- Annual mean forcing
- **Mixed boundary conditions**

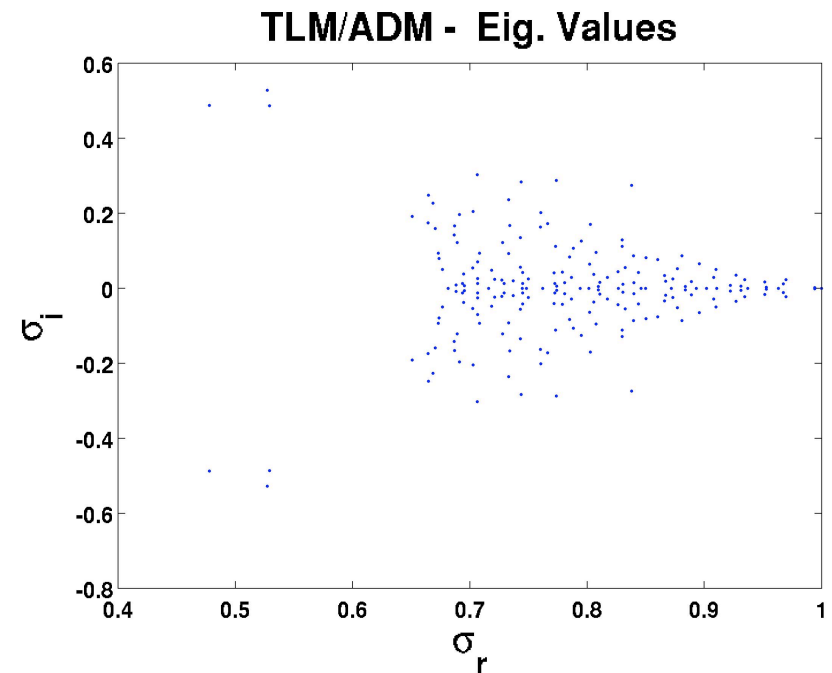
(e.g., Marshall et al. 1997; <http://mitgcm.org>)



Stability of the Tangent Linear Model



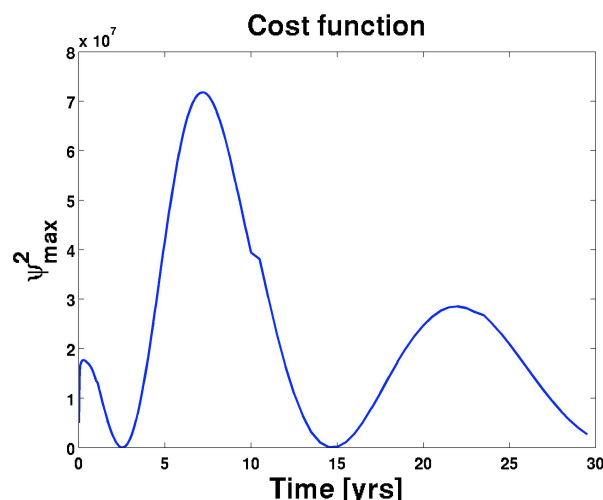
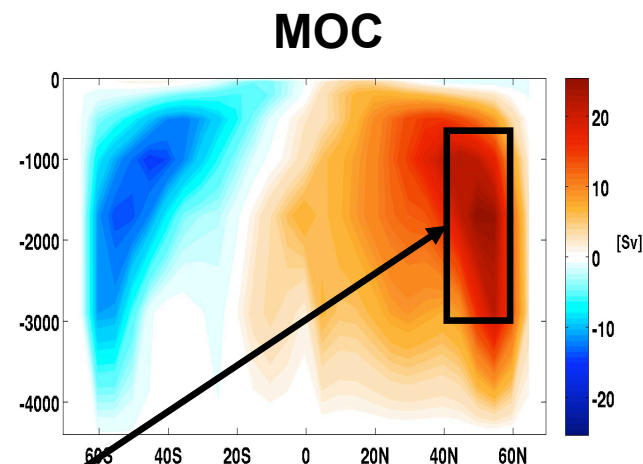
TLM least damped mode with decay time of 800 yrs



TLM/ADM Imag. vs Real eig. values for $t=2$ yrs

Transient growth of MOC anomalies: preliminary results *(Zanna et al, in prep)*

- **MOC stability & variability:** salinity advective feedback, from interannual to multidecadal, NAO-gyre interaction (*e.g, Marotzke 1990, Marshall et al 2001*)
- Few studies on transient amplification of MOC (*e.g., Lohman & Schneider 1999, Zanna & Tziperman 2005, Sevellec et al 2008*)

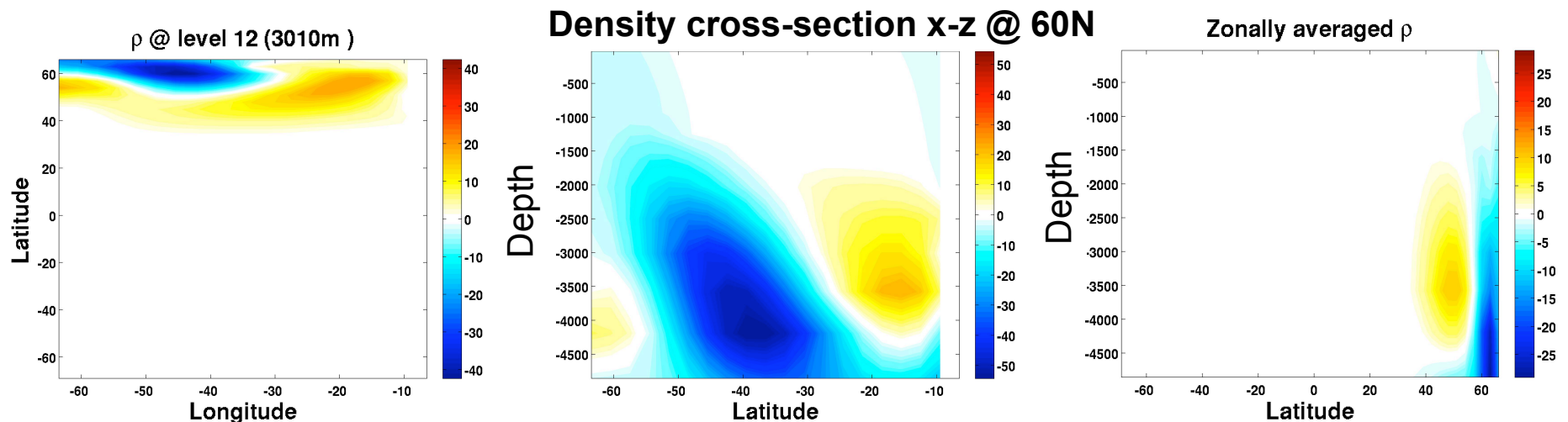


Cost function as fct of time when initializing TLM with optimals

- **Cost function:** sum of the square of the MOC anomalies 50N-60N, 1000m-3500m (*Bugnion & Hill 2006*)
- **Results:** Optimal i.c of T & S lead to growth of MOC anomalies after ~8 yrs associated with non-orthogonal oscillatory modes. Growth factor ~8.

Transient growth of MOC anomalies: Initial conditions

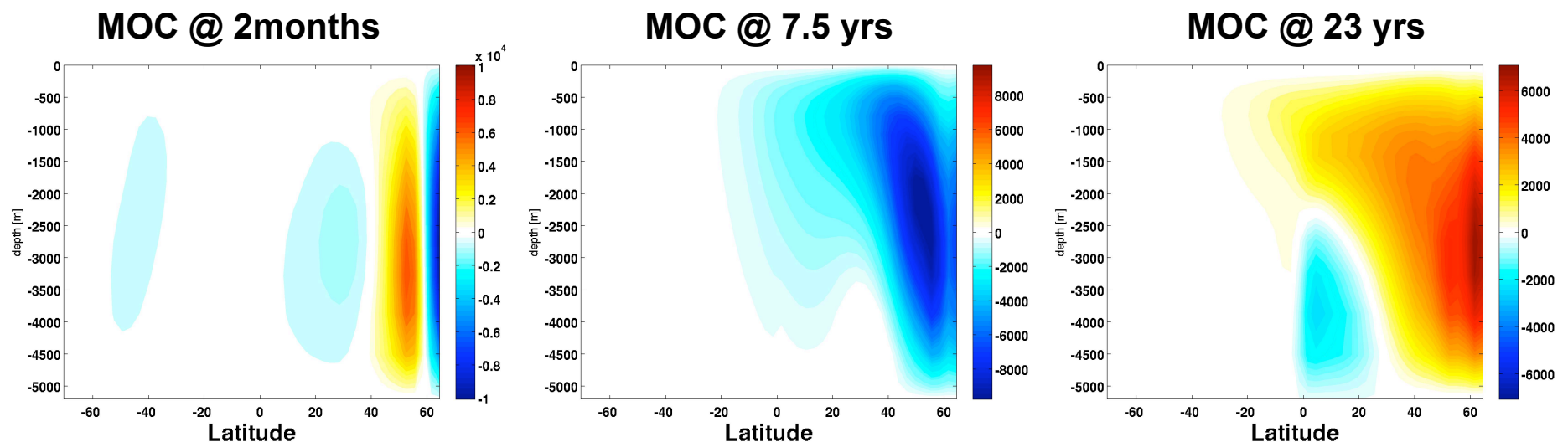
- Signal mostly in NH with **baroclinic structure**
- **Strong signal in the deep ocean** with additive contribution of T & S to buoyancy
- T & S necessary for growth (*unlike Marotzke 1990; Sevellec et al 2008*)
- Similarities with unstable oscillatory mode under fixed flux of *Raa & Dijkstra (2002)*



Growth Mechanism:

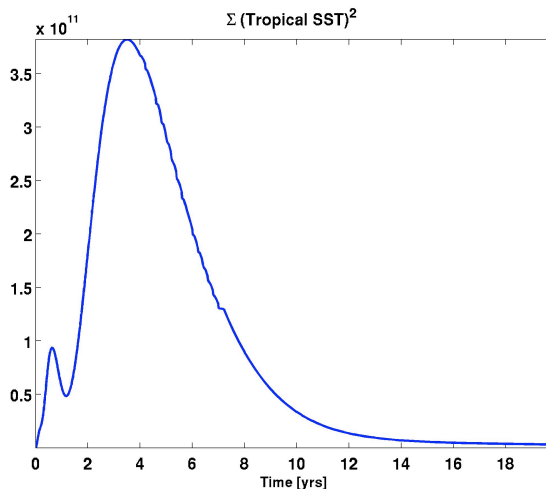
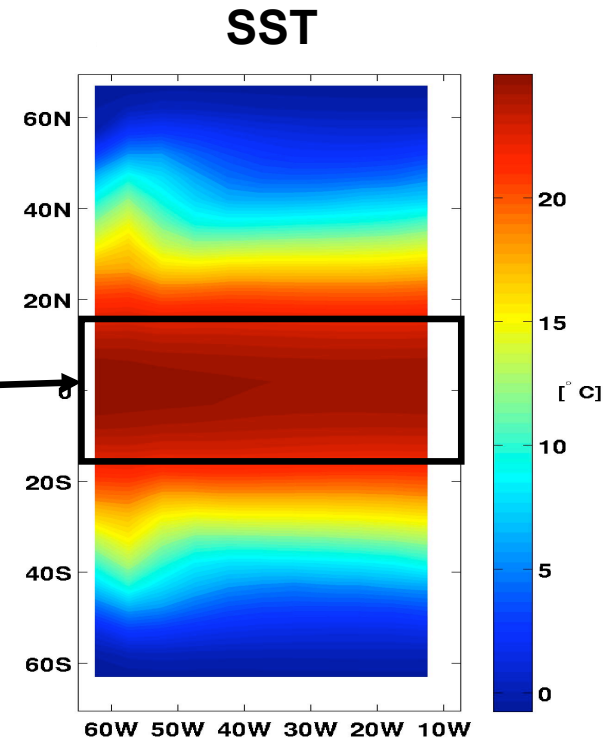
Preliminary & simplified results

- **Several decaying oscillatory modes** under mixed boundary conditions but no growth of individual modes
- **Growth** = change of APE due advection of density perturbations by the mean flow creating N-S & E-W buoyancy gradient
- **Oscillation** = phase difference btw the N-S & E-W buoyancy gradients



Exciting Tropical SST anomalies

- Tropical Atlantic Variability **mechanisms** → **air-sea interaction** or connected to **seasonal cycle** (e.g., Chang, Xie & Carton, Jochum et al)
- **Cost function**: sum of the square of the SST anomalies btw 15S & 15 N
- **Optimal i.c.** = deep salinity anomalies near the western boundary @ 30N/S & 50 N/S



Sum of squares of Tropical SST anomalies as fct of time

- **Results**: Optimal growth of tropical SST anomalies after 3.5-4 yrs. Initial 0.1 ppt → 0.4 C.
- **Mechanism**: geostrophic adjustment → Coastal & Equatorial Kelvin waves (Zanna et al, submitted to JPO)

Conclusions

- *Small* perturbations → *Large* amplification on interannual timescales without unstable modes
- **Identification of new mechanisms** leading to growth of perturbations
- Preferred anomalies located in the **deep ocean**:
- **Non-normal dynamics can possibly play a dominant role in generating variability on interannual timescales if excited by stochastic forcing**

...

Conclusions

- From Box models to GCMs:
 - non-normality of the propagator mainly due to advection & surface boundary conditions
 - Faster time scales in 3D (<10yrs) than in 2D models (decades)
 - More complex dynamics in full GCM & possibility to explore different physical quantities: MOC, energy, heat flux, etc
- Idealized MITgcm:
 - Eigenvectors of the TLM &
 - Singular vectors for different physical quantities (in //)
- Challenges:
 - Calculations are relatively expensive for higher resolutions
 - Bathymetry: strong sensitivity in shallow areas
 - Atmosphere (non-normality increases when introducing atmospheric coupling)